

Operators

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Operators: When an operator meets a ket, it produces another ket.
(cf. bra: When a bra meets a ket, it produces a scalar.)

- $\hat{M}|A\rangle = |B\rangle.$

Linear Operators

- $\hat{M}z|A\rangle = z|B\rangle.$
- $\hat{M}(|A\rangle + |B\rangle) = \hat{M}|A\rangle + \hat{M}|B\rangle.$

Matrix Representations

We have two kets and a operator with the following relation:

$|A\rangle = \sum_i a_i |i\rangle$, $|B\rangle = \sum_i b_i |i\rangle$, and

$$\hat{M}|A\rangle = |B\rangle.$$

Then,

$$\sum_i \hat{M} a_i |i\rangle = \sum_i \hat{M} |i\rangle a_i = \sum_i b_i |i\rangle.$$

Applying another basis bra $\langle j|$,

$$\sum_i \langle j | \hat{M} | i \rangle a_i = \sum_i b_i \langle j | i \rangle = \sum_i b_i \delta_{ij} = b_j. \quad (1)$$

Matrix Representations

$$\sum_i \langle j | \hat{M} | i \rangle a_i = b_j.$$

If we represent $\langle j | \hat{M} | i \rangle$ as an element of a matrix,

$$\begin{pmatrix} \langle 1 | \hat{M} | 1 \rangle & \langle 1 | \hat{M} | 2 \rangle & \langle 1 | \hat{M} | 3 \rangle & \cdots \\ \langle 2 | \hat{M} | 1 \rangle & \langle 2 | \hat{M} | 2 \rangle & \langle 2 | \hat{M} | 3 \rangle & \cdots \\ \langle 3 | \hat{M} | 1 \rangle & \langle 3 | \hat{M} | 2 \rangle & \langle 3 | \hat{M} | 3 \rangle & \cdots \\ \vdots & & & \end{pmatrix} \begin{pmatrix} a_1 \\ a_2 \\ a_3 \\ \vdots \end{pmatrix} = \begin{pmatrix} b_1 \\ b_2 \\ b_3 \\ \vdots \end{pmatrix}, \quad (2)$$

then the meaning of this operator is transparent.

- Now we accept that as an axiom, all physical observable in quantum mechanics are operators such as \hat{x} , \hat{p} , \hat{H} , and \hat{L} .
- In general two matrices do not commute $[A, B] \neq 0$.
- Observable must be real: Hermitian operators.

Hermitian Operators

Observable must be “real” in the “real” world.

- $\langle \Psi | \hat{M} \Psi \rangle = \langle \hat{M} \Psi | \Psi \rangle.$

In quantum mechanics, observable must be Hermitian operators.

- $\langle f | \hat{M} g \rangle = \langle \hat{M} f | g \rangle$ for all f and g .
- $\hat{M} = \hat{M}^\dagger = [\hat{M}^T]^*$
- $m_{ji} = m_{ij}^*$
- Transpose (T) and complex conjugate ($*$)



Charles Hermite
(1822 ~ 1901)

Eigenvalues and Eigenvectors

Eigenvectors $|\lambda\rangle$ are special vectors that are transformed into scalar multiples of themselves after linear transformation,

$$\hat{M}|\lambda\rangle = \lambda|\lambda\rangle, \quad (3)$$

where the complex number λ is their eigenvalue. Discriminate between λ and $|\lambda\rangle$.

Eigenvalues and Eigenvectors

How to obtain eigenvectors and eigenvalues?

1. Start from $\hat{M}|\lambda\rangle = \lambda|\lambda\rangle$
2. $(\hat{M} - \lambda\mathbf{I})|\lambda\rangle = |0\rangle$
3. Assuming $|\lambda\rangle$ is nonzero, solve $\det(\hat{M} - \lambda\mathbf{I})$.
4. Solving the characteristic equation, we obtain eigenvalues of \hat{M} .
5. By putting eigenvalues back into the eigenvector equation, we can finally find eigenvectors.

Yes. You have to solve many practice problems of eigenvector equations.

Eigenvalues and Eigenvectors in Quantum Mechanics

Time-independent Schrödinger equation:

$$\hat{H}\psi = E\psi. \quad (4)$$

Thus, in a perspective of linear algebra

- \hat{H} : operator
- E_n : eigenvalues
- ψ_n : eigenvectors

Hermitian Operators

- Eigenvalues of a Hermitian operator are always real. They are same as their complex conjugate.
- Two different eigenvectors corresponding two different eigenvalues are orthogonal.
- When two eigenvalues are the same (degeneracy), we can choose two orthogonal eigenvectors (Gram-Schmidt orthogonalization).
- Eigenvectors of a Hermitian operator are a complete set. Any vectors can be decomposed by the sum of eigenvectors.

Real eigenvalues

$$\begin{aligned}\hat{M}|\lambda\rangle &= \lambda|\lambda\rangle, \\ \langle\lambda|\hat{M}^\dagger &= \langle\lambda|\lambda^*.\end{aligned}$$

Suppose that \hat{M} is a Hermitian operator ($\hat{M} = \hat{M}^\dagger$),

$$\begin{aligned}\hat{M}|\lambda\rangle &= \lambda|\lambda\rangle, \\ \langle\lambda|\hat{M} &= \langle\lambda|\lambda^*.\end{aligned}$$

Then, we have (applying $\langle\lambda|$ bra or $|\lambda\rangle$ ket),

$$\begin{aligned}\langle\lambda|\hat{M}|\lambda\rangle &= \lambda\langle\lambda|\lambda\rangle, \\ \langle\lambda|\hat{M}|\lambda\rangle &= \lambda^*\langle\lambda|\lambda\rangle.\end{aligned}$$

Therefore the eigenvalues of a Hermitian operator are real ($\lambda = \lambda^*$).

Orthogonal Eigenvectors

$$\hat{M}|\lambda_1\rangle = \lambda_1|\lambda_1\rangle,$$

$$\hat{M}|\lambda_2\rangle = \lambda_2|\lambda_2\rangle,$$

Since \hat{M} is a Hermitian operator,

$$\langle\lambda_1|\hat{M} = \lambda_1\langle\lambda_1|,$$

$$\hat{M}|\lambda_2\rangle = \lambda_2|\lambda_2\rangle,$$

Then, we have (applying $\langle\lambda_1|$ bra or $|\lambda_2\rangle$ ket),

$$\langle\lambda_1|\hat{M}|\lambda_2\rangle = \lambda_1\langle\lambda_1|\lambda_2\rangle,$$

$$\langle\lambda_1|\hat{M}|\lambda_2\rangle = \lambda_2\langle\lambda_1|\lambda_2\rangle.$$

Finally, we have

$$(\lambda_1 - \lambda_2)\langle\lambda_1|\lambda_2\rangle = 0.$$

Degeneracy and Gram-Schmidt Orthogonalization

If $\lambda_1 = \lambda_2$, we can choose two orthogonal eigenvectors by using Gram-Schmidt orthogonalization. You will prove it (assignment 2).

Eigenvectors Decomposition

When we introduce a set of orthonormal eigenvectors $|i\rangle$, any ket $|A\rangle$ can be represented by a linear combination of the eigenvectors:

$$|A\rangle = \sum_i \alpha_i |i\rangle. \quad (5)$$

Where to go next...

Unitary Operators and Quantum Dynamics.