

Quantum Mechanics in Three Dimensions

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Schrödinger Equation in Three Dimensions

Schrödinger equation for the Hamiltonian operator \hat{H} in three dimensions is still given by

$$i\hbar \frac{\partial \Psi}{\partial t} = \hat{H} \Psi,$$

where the Hamiltonian is

$$\hat{H} = \frac{1}{2m} (p_x^2 + p_y^2 + p_z^2) + V.$$

By the standard prescription

$$\hat{p} \rightarrow -i\hbar \nabla, \quad (\hat{p}_i \rightarrow -i\hbar \frac{\partial}{\partial i})$$

we finally obtain

$$i\hbar \frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \Psi + V \Psi. \quad (1)$$

General Solution to the Schrödinger Equation

When the potential V is independent of time, the general solution to the Schrödinger equation in three dimensions is

$$\Psi(\vec{r}, t) = \sum_n c_n \psi_n(\vec{r}) e^{-iE_n t/\hbar},$$

where ψ_n is the eigenfunctions with E_n of the time-independent Schrödinger equation:

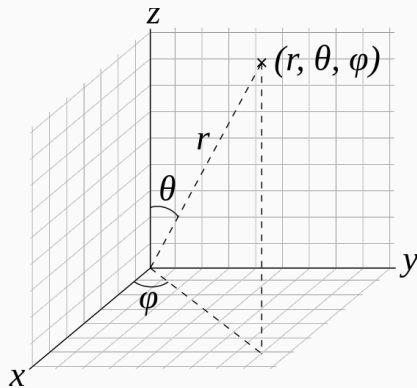
$$-\frac{\hbar^2}{2m} \nabla^2 \psi + V \psi = E \psi.$$

Note that the canonical commutation relations satisfy [see Problem (4.1) in Griffith]:

$$[r_i, p_j] = i\hbar \delta_{ij}, \quad [r_i, r_j] = [p_i, p_j] = 0.$$

Schrödinger Equation in Spherical Coordinates

$$x = r \sin \theta \cos \varphi, \quad y = r \sin \theta \sin \varphi, \quad z = r \cos \theta.$$



Spherical coordinate system (r, θ, φ)

Spherical Coordinate System

$$x = r \sin \theta \cos \varphi, \quad y = r \sin \theta \sin \varphi, \quad z = r \cos \theta.$$

Scale factor:

$$ds_i = h_i dq_i, \quad ds^2 = \sum_i (h_i dq_i)^2,$$

where

$$h_i^2 = \frac{\partial x}{\partial q_i} \frac{\partial x}{\partial q_i} + \frac{\partial y}{\partial q_i} \frac{\partial y}{\partial q_i} + \frac{\partial z}{\partial q_i} \frac{\partial z}{\partial q_i}.$$

In Spherical Coordinate System,

$$h_r = 1, \quad h_\theta = r, \quad h_\varphi = r \sin \theta.$$

Lapalcan operator in Spherical Coordinate System

$$h_r = 1, \quad h_\theta = r, \quad h_\varphi = r \sin \theta.$$

$$\nabla\psi = \frac{1}{h_1} \frac{\partial\psi}{\partial q_1} \hat{q}_1 + \frac{1}{h_2} \frac{\partial\psi}{\partial q_2} \hat{q}_2 + \frac{1}{h_3} \frac{\partial\psi}{\partial q_3} \hat{q}_3,$$

$$\nabla \cdot V = \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial q_1} (V_1 h_2 h_3) + \frac{\partial}{\partial q_2} (V_2 h_3 h_1) + \frac{\partial}{\partial q_3} (V_3 h_1 h_2) \right],$$

$$\nabla^2\psi = \nabla \cdot \nabla\psi$$

$$= \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial q_1} \left(\frac{h_2 h_3}{h_1} \frac{\partial\psi}{\partial q_1} \right) + \frac{\partial}{\partial q_2} \left(\frac{h_3 h_1}{h_2} \frac{\partial\psi}{\partial q_2} \right) + \frac{\partial}{\partial q_3} \left(\frac{h_1 h_2}{h_3} \frac{\partial\psi}{\partial q_3} \right) \right].$$

Thus, in spherical coordinates

$$\nabla^2 = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \phi^2}.$$

Schrödinger Equation in Spherical Coordinates

$$-\frac{\hbar^2}{2m} \left[\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \left(\frac{\partial^2 \psi}{\partial \phi^2} \right) \right] + V\psi = E\psi.$$

Then, we attempt the method of separation of variables for $V = V(r)$,

$$\psi(r, \theta, \phi) = R(r)Y(\theta, \phi).$$

Then, we have radial and angular equations:

$$\frac{1}{R} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) - \frac{2mr^2}{\hbar^2} [V(r) - E] = l(l+1),$$
$$\frac{1}{Y} \left[\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial Y}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2 Y}{\partial \phi^2} \right] = -l(l+1).$$

Radial Equation

The radial equation of Schrödinger equation:

$$\frac{1}{R} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) - \frac{2mr^2}{\hbar^2} [V(r) - E] = l(l+1).$$

Introducing a variable $u = rR(r)$, we find

$$-\frac{\hbar^2}{2m} \frac{d^2u}{dr^2} + \left[V + \frac{\hbar^2}{2m} \frac{l(l+1)}{r^2} \right] u = Eu.$$

If we define the effective potential,

$$V_e = V + \frac{\hbar^2}{2m} \frac{l(l+1)}{r^2},$$

we find the equation that is identical to the one-dimensional Schrödinger equation:

$$-\frac{\hbar^2}{2m} \frac{d^2u}{dr^2} + V_e u = Eu. \quad (2)$$

Angular Equation

The angular equation of Schrödinger equation:

$$\frac{1}{Y} \left[\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial Y}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2 Y}{\partial \phi^2} \right] = -l(l+1).$$

The general solution of Y is called by spherical harmonics:

$$Y_l^m(\theta, \phi) = \epsilon \sqrt{\frac{(2l+1)(l-|m|)!}{4\pi(l+|m|)!}} e^{im\phi} P_l^m(\cos \theta), \quad (3)$$

where $\epsilon = (-1)^m$ for $m \geq 0$ and $\epsilon = 1$ for $m \leq 0$ and P_l^m is the associated Legendre function.