

# Metals and the Fermi Gas

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# Electrons in a Metal

Consider electrons in a metal as free non-interacting electrons.

$$\begin{aligned}\langle n \rangle &= \frac{1}{e^{\beta(\epsilon - \mu)} + 1}, \\ \epsilon &= \frac{\hbar^2 k^2}{2m}, \\ g(\epsilon) d\epsilon &= \frac{mV}{\pi^2 \hbar^3} \sqrt{2m\epsilon} d\epsilon.\end{aligned}$$

# Electrons in a Metal

Then, the total number of electrons is

$$N = \int_0^{\infty} g(\epsilon) \langle n \rangle d\epsilon,$$
$$\langle n \rangle = \frac{1}{e^{\beta(\epsilon-\mu)} + 1}.$$

At  $T = 0$  (equivalently  $\beta \rightarrow \infty$ ),

- When  $\epsilon \leq \mu$ ,  $\langle n \rangle = 1$ .
- When  $\epsilon > \mu$ ,  $\langle n \rangle = 0$ .

# Fermi Energy

We define the zero temperature value of chemical potential as Fermi energy  $\epsilon_F$ . It plays an important role to analyze the (electrical) property of solids. And, we can obtain the total number of particles by using  $\epsilon_F$  as

$$\begin{aligned}
 N &= \int_0^{\infty} g(\epsilon) \langle n \rangle d\epsilon \\
 &= \int_0^{\epsilon_F} g(\epsilon) d\epsilon \\
 &= \int_0^{\epsilon_F} \frac{mV}{\pi^2 \hbar^3} \sqrt{2m\epsilon} d\epsilon \\
 &= \frac{(2\epsilon_F m)^{3/2}}{3\pi^2 \hbar^3} V.
 \end{aligned}$$

This approximation is not accurate in general because we assume independent free particles, but it is surprisingly useful as the simplest example.

# Degeneracy Pressure

The Fermi energy is then

$$\epsilon_F = \frac{\hbar^2}{2m} (3\pi^2 N/V)^{2/3} \sim \left( \frac{N}{V} \right)^{2/3}.$$

The energy is

$$\begin{aligned} E &= \int_0^{\epsilon_F} \epsilon g(\epsilon) d\epsilon = \int_0^{\epsilon_F} \epsilon \frac{mV}{\pi^2 \hbar^3} \sqrt{2m\epsilon} d\epsilon \\ &\sim V \epsilon_F^{5/2} = V \left( \frac{N}{V} \right)^{5/3}. \end{aligned}$$

The degeneracy pressure is

$$P \sim \frac{E}{V} = \left( \frac{N}{V} \right)^{5/3} = \rho^{5/3}.$$